

DIFFERENTIAL OPERATORS ON BUNDLE OF ACCELERATIONS IN HOMOGENEOUS APPROACH

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Abstract. Using a special (h, v_1, v_2) metrical structure on the bundle of accelerations, the divergence of a field, the gradient and the Laplacean of a function on $Osc^{(2)}M$ are defined in this paper. In the case of homogeneous lift of a Riemannian metric to bundle of accelerations we obtain some results on differential operators.

MSC: 53C60

Keywords and phrases: bundle of accelerations, (α, β, η) Sasaki lift to bundle of accelerations of a Riemannian structure, homogeneous Sasaki lift, differential operator.

1. Introduction

Let $(E = Osc^{(2)}M, \pi^2, M)$ be the 2-osculator bundle of a real, smooth, n -dimensional manifold M . For a local chart $(U, \varphi = (x^i))$ on M , its induced local chart on $E = Osc^{(2)}M$ will be denoted by $((\pi^2)^{-1}(U), \Phi = (x, y^{(1)i}, y^{(2)i}))$.

Let $R^{(n)} = (M, \gamma)$ be a Riemannian space, M being a real n -dimensional differentiable manifold and γ a Riemannian metric on M , having the local coordinates $\gamma_{ij}(x)$, $x \in U \subset M$. We can extend γ_{ij} to $\pi^{-1}(U) \subset E$, setting:

$$(\gamma_{ij} \circ \pi)(u) = \gamma_{ij}(x), u \in \pi^{-1}(U), \pi(u) = x \quad (1.1.)$$

In this case $\gamma_{ij} \circ \pi$ gives a d-tensor field on E . We also denoted it by γ_{ij} . If we denote by γ_{jk}^i the Christoffel symbols of γ we have the following result:

Theorem 1.1. *The coefficients of the nonlinear connection on \tilde{E} determined only by the Riemannian structure $\gamma(x)$ are given by:*

$$\begin{cases} N^i_{(1)j}(x, y^{(1)}) = \gamma^i_{j0} \\ N^i_{(2)j}(x, y^{(1)}, y^{(2)}) = \frac{1}{2}A^i_j + \gamma^i_{j\bar{0}} \end{cases} \quad (1.2)$$

where "0" means the contraction by $(y^{(1)})$, " $\bar{0}$ " means the contraction by $(y^{(2)})$ and $A^i_j = \left(\frac{\partial \gamma^i_{j0}}{\partial x^p} y^{(1)p} - \gamma^i_{0m} \cdot \gamma^m_{j0} \right)$.

The nonlinear connection N on \tilde{E} provides the existence of the adapted basis $(\delta_k, \delta^{(1)}_k, \delta^{(2)}_k)$ of the tangent space $T_u E$. These vector fields are given by the following relation:

$$\begin{cases} \delta_k = \frac{\partial}{\partial x^k} - N^i_{(1)k} \cdot \frac{\partial}{\partial y^{(1)i}} - N^i_{(2)k} \cdot \frac{\partial}{\partial y^{(2)i}} \\ \delta^{(1)}_k = \frac{\partial}{\partial y^{(1)k}} - N^i_{(1)k} \cdot \frac{\partial}{\partial y^{(1)i}} \\ \delta^{(2)}_k = \frac{\partial}{\partial y^{(2)k}} \end{cases} \quad (1.3)$$

Further, we need the following result with respect to the vector fields of the adapted basis $(\delta_k, \delta^{(1)}_k, \delta^{(2)}_k)$:

Theorem 1.2. *In the case of the nonlinear connection on \tilde{E} determined only by the Riemannian structure $\gamma(x)$, the local coefficients of Lie brackets of the vector fields of the adapted basis $(\delta_k, \delta^{(1)}_k, \delta^{(2)}_k)$ are given by:*

$$\begin{cases} R^i_{(01)jk} = r^i_{0jk} & R^i_{(12)jk} = 0 \end{cases} \quad (1.4)$$

$$\begin{cases} R^i_{(02)jk} = \frac{1}{2} \frac{\partial r^i_{0jk}}{\partial x^p} \cdot y^{(1)p} + \\ \frac{1}{2} \cdot (\gamma^i_{0m} \cdot r^m_{0jk} + \gamma^m_{0j} \cdot r^i_{0km} + \gamma^m_{0k} \cdot r^i_{0mj}) + r^i_{0jk} \end{cases} \quad (1.5)$$

$$\begin{cases} B^i_{(11)jk} = \gamma^i_{jk} & B^i_{(22)jk} = \gamma^i_{jk} \\ B^i_{(12)jk} = \frac{1}{2} \cdot r^i_{0jk} + \gamma^i_{0m} \cdot \gamma^m_{jk} & B^i_{(21)jk} = 0 \end{cases} \quad (1.6)$$

where $r_{h,jk}^i(x)$ are the local coefficients of the curvature tensor field of the Riemannian space $(M, \gamma_{ij}(x))$.

Theorem 1.3. ([6]) The pair $\text{Pr ol}^2 R^{(n)} = (\tilde{Osc}^{(2)} M, G)$, where:

$$G = \gamma_{ij}(x) \cdot dx^i \otimes dx^j + \gamma_{ij}(x) \cdot \delta y^{(1)i} \otimes \delta y^{(1)j} + \gamma_{ij}(x) \cdot \delta y^{(2)i} \otimes \delta y^{(2)j} \quad (1.7.)$$

is a Riemannian space of dimension $3n$, whose metric structure G depends only on the structure $\gamma(x)$ of the a priori given Riemann space $R^{(n)} = (M, \gamma)$.

We say that G is Sasaki N - lift of the Riemannian structure γ .

Next, let us consider the homothety $h_t : (x, y^{(1)}, y^{(2)}) \rightarrow (x, ty^{(1)}, t^2 y^{(2)})$, $t \in R^*$ on the fibres of fibred bundle $Osc^{(2)} M$. We notice that G is transformed as follows:

$$G \circ h_t(x, y^{(1)}, y^{(2)}) = \gamma_{ij}(x) \cdot dx^i \otimes dx^j + t^2 \cdot \gamma_{ij}(x) \cdot \delta y^{(1)i} \otimes \delta y^{(1)j} + t^4 \cdot \gamma_{ij}(x) \cdot \delta y^{(2)i} \otimes \delta y^{(2)j} \quad (1.8.)$$

Therefore we have:

Theorem 1.4. The Sasaki lift G is not homogeneous on the fibres of fibres bundles $Osc^{(2)} M$.

2. The (α, β, η) Sasaki lift to bundle of accelerations of a Riemannian metric

Definition 2.1. We call the $(\alpha\beta\eta)$ lift to $\tilde{Osc}^{(2)} M$ of the fundamental tensor field γ_{ij} of a Riemannian space $R^{(n)}$ the following tensor field on $\tilde{Osc}^{(2)} M$:

$$\overset{(\alpha\beta\eta)}{G} = g_{ij}(x, y^{(1)}) \cdot dx^i \otimes dx^j + h_{ij}(x, y^{(1)}) \cdot \delta y^{(1)i} \otimes \delta y^{(1)j} + l_{ij}(x, y^{(1)}) \cdot \delta y^{(2)i} \otimes \delta y^{(2)j} \quad (2.1.)$$

where we put:

$$\begin{cases} g_{ij}(x, y^{(1)}) = \alpha(F^2) \cdot \gamma_{ij}(x) \\ h_{ij}(x, y^{(1)}) = \beta(F^2) \cdot \gamma_{ij}(x) \\ l_{ij}(x, y^{(1)}) = \eta(F^2) \cdot \gamma_{ij}(x) \end{cases} \quad (2.2.)$$

where F is the Finsler function of Riemannian space $R^{(n)}$: $F^2 = \gamma_{ij}(x) \cdot y^{(1)i} \cdot y^{(1)j}$, $y_i^{(1)} = \gamma_{ij}(x) \cdot y^{(1)j}$ and $\alpha, \beta, \eta: R_+ \rightarrow R_+$ are differentiable functions.

Theorem 2.1. The following properties hold:

- The pair $(\tilde{Osc}^{(2)} M, \tilde{G}^{(\alpha\beta\eta)})$ is a Riemannian space;
- $\tilde{G}^{(\alpha\beta\eta)}$ depends only on the Riemannian metric γ of the Riemann space $R^{(n)}$;
- The distributions N, V_1, V_2 are orthogonal with respect to $\tilde{G}^{(\alpha\beta\eta)}$.

Remarks:

- For $\alpha = \beta = \eta = 1$ we obtain the Sasaki lift introduced by R. Miron in the monograph [2].
- For $\alpha(t) = 1, \beta(t) = \frac{a^2}{t}, \eta(t) = \frac{a^4}{t^2}$ ($a > 0$) we obtain a new lift called homogeneous Sasaki lift, which will be denoted with $\tilde{G}^{(0)}$. Therefore we have:

$$\begin{aligned} \tilde{G}^{(0)} = & \gamma_{ij}(x) \cdot dx^i \otimes dx^j + \frac{a^2}{F^2} \cdot \gamma_{ij}(x) \cdot \delta y^{(1)i} \otimes \delta y^{(1)j} \\ & + \frac{a^4}{F^4} \cdot \gamma_{ij}(x) \cdot \delta y^{(2)i} \otimes \delta y^{(2)j} \end{aligned} \quad (2.3.)$$

Definition 2.2. A linear connection D on $\tilde{Osc}^{(2)} M$ is called a $(\alpha\beta\eta)$ metrical connection if $D \tilde{G}^{(\alpha\beta\eta)} = 0$ and D preserves by parallelism the horizontal distribution N .

In the adapted basis, any linear connection on E can be represented in the following way:

$$\begin{cases} D_{\delta_k} \delta_j = L_{jk}^{(H)} \cdot \delta_i + L_{jk}^{(1)} \cdot \delta_i + L_{jk}^{(2)} \cdot \delta_i \\ D_{\delta_k}^{(1)} \delta_j = L_{jk}^{(1)} \cdot \delta_i + L_{jk}^{(V_1)} \cdot \delta_i + L_{jk}^{(4)} \cdot \delta_i \\ D_{\delta_k}^{(2)} \delta_j = L_{jk}^{(2)} \cdot \delta_i + L_{jk}^{(6)} \cdot \delta_i + L_{jk}^{(V_2)} \cdot \delta_i \end{cases} \quad (2.4.1)$$

$$\begin{cases} D_{\delta_k}^{(1)} \delta_j = F_{jk}^{(H)^i} \cdot \delta_i + F_{jk}^{(1)^i} \cdot \delta_i + F_{jk}^{(2)^i} \cdot \delta_i \\ D_{\delta_k}^{(1)} \delta_j = F_{jk}^{(1)^i} \cdot \delta_i + F_{jk}^{(v_1)^i} \cdot \delta_i + F_{jk}^{(4)^i} \cdot \delta_i \\ D_{\delta_k}^{(1)} \delta_j = F_{jk}^{(2)^i} \cdot \delta_i + F_{jk}^{(5)^i} \cdot \delta_i + F_{jk}^{(v_2)^i} \cdot \delta_i \end{cases} \quad (2.4.2)$$

$$\begin{cases} D_{\delta_k}^{(2)} \delta_j = C_{jk}^{(H)^i} \cdot \delta_i + C_{jk}^{(1)^i} \cdot \delta_i + C_{jk}^{(2)^i} \cdot \delta_i \\ D_{\delta_k}^{(2)} \delta_j = C_{jk}^{(1)^i} \cdot \delta_i + C_{jk}^{(v_1)^i} \cdot \delta_i + C_{jk}^{(4)^i} \cdot \delta_i \\ D_{\delta_k}^{(2)} \delta_j = C_{jk}^{(2)^i} \cdot \delta_i + C_{jk}^{(5)^i} \cdot \delta_i + C_{jk}^{(v_2)^i} \cdot \delta_i \end{cases} \quad (2.4.3)$$

The systems of functions $(L_{jk}^{(H)^i}, \dots, C_{jk}^{(v_2)^i})$ are the coefficients of D . We can prove the following result:

Theorem 3.3. *There exist the $(\alpha\beta\eta)$ metrical connections on $\widetilde{OSC}^{(2)} M$, which depend on the metric tensor γ .*

One of them has the following coefficients:

$$L_{jk}^{(1)^i} = L_{jk}^{(2)^i} = L_{jk}^{(3)^i} = L_{jk}^{(4)^i} = L_{jk}^{(5)^i} = L_{jk}^{(6)^i} = 0 \quad (2.5.)$$

$$F_{jk}^{(1)^i} = F_{jk}^{(2)^i} = F_{jk}^{(3)^i} = F_{jk}^{(4)^i} = F_{jk}^{(5)^i} = F_{jk}^{(6)^i} = 0 \quad (2.6.)$$

$$C_{jk}^{(H)^i} = C_{jk}^{(1)^i} = C_{jk}^{(2)^i} = C_{jk}^{(3)^i} = C_{jk}^{(4)^i} = C_{jk}^{(5)^i} = C_{jk}^{(6)^i} = C_{jk}^{(v_1)^i} = C_{jk}^{(v_2)^i} = 0 \quad (2.7.)$$

$$L_{jk}^{(H)^i} = L_{jk}^{(v_1)^i} = L_{jk}^{(v_2)^i} = \gamma_{jk}^i \quad (2.8.)$$

$${}^{(H)}F_{jk}^i = \frac{\alpha'}{\alpha} \cdot \Lambda_{jk}^i, \quad {}^{(V_1)}F_{jk}^i = \frac{\beta'}{\beta} \cdot \Lambda_{jk}^i, \quad {}^{(V_2)}F_{jk}^i = \frac{\eta'}{\eta} \cdot \Lambda_{jk}^i \quad (2.9)$$

where:

$$\Lambda_{jk}^i = \delta_j^i y_k^{(1)} + \delta_k^i y_j^{(1)} - \gamma_{jk} y^{(1)i} \quad (2.10)$$

The simplicity of this metrical N-connection and the fact that it is determined only by the Riemannian metric γ , allows to call it the *canonical* $(\alpha\beta\eta)$ metrical connection of the space $(\tilde{T}Osc^{(2)} M, G^{(\alpha\beta\eta)})$. In the adapted basis $(\delta_k^{(1)}, \delta_k^{(2)}, \delta_k^{(3)})$ the torsion tensor field of the canonical $(\alpha\beta\eta)$ metrical connection of the space $(\tilde{T}Osc^{(2)} M, G^{(\alpha\beta\eta)})$ satisfies the following relations:

$$\left\{ \begin{array}{l} T(\delta_k, \delta_j) = T_{jk}^{(1)} \cdot \delta_i + T_{jk}^{(2)} \cdot \delta_i + T_{jk}^{(3)} \cdot \delta_i \\ T(\delta_k^{(1)}, \delta_j) = P_{jk}^{(1)} \cdot \delta_i + P_{jk}^{(2)} \cdot \delta_i + P_{jk}^{(3)} \cdot \delta_i \\ T(\delta_k^{(2)}, \delta_j) = Q_{jk}^{(1)} \cdot \delta_i + Q_{jk}^{(2)} \cdot \delta_i + Q_{jk}^{(3)} \cdot \delta_i \\ T(\delta_k^{(1)}, \delta_j^{(1)}) = S_{jk}^{(1)} \cdot \delta_i + S_{jk}^{(2)} \cdot \delta_i + S_{jk}^{(3)} \cdot \delta_i \\ T(\delta_k^{(2)}, \delta_j^{(1)}) = V_{jk}^{(1)} \cdot \delta_i + V_{jk}^{(2)} \cdot \delta_i + V_{jk}^{(3)} \cdot \delta_i \\ T(\delta_k^{(2)}, \delta_j^{(2)}) = U_{jk}^{(1)} \cdot \delta_i + U_{jk}^{(2)} \cdot \delta_i + U_{jk}^{(3)} \cdot \delta_i \end{array} \right. \quad (2.11)$$

Theorem 2.4. The local components of the torsion tensor field of the canonical $(\alpha\beta\eta)$ metrical connection are given by:

$$T_{jk}^{(1)} = P_{jk}^{(2)} = Q_{jk}^{(1)} = Q_{jk}^{(2)} = Q_{jk}^{(3)} = 0 \quad (2.12)$$

$$S_{jk}^{(1)} = S_{jk}^{(2)} = S_{jk}^{(3)} = V_{jk}^{(1)} = V_{jk}^{(2)} = U_{jk}^{(1)} = U_{jk}^{(2)} = U_{jk}^{(3)} = 0 \quad (2.13)$$

$$T_{jk}^{(2)} = R_{(01)jk}^{(2)}, \quad T_{jk}^{(3)} = R_{(02)jk}^{(3)}, \quad (2.14.)$$

$$P_{jk}^{(1)} = F_{jk}^{(H)}, \quad P_{jk}^{(3)} = B_{(12)jk}^{(3)}, \quad V_{jk}^{(3)} = -F_{jk}^{(V_3)} \quad (2.15.)$$

3. Differential operators

Let $F(\tilde{E})$ be the ring of the smooth functions on \tilde{E} and let $X(\tilde{E})$ be the $F(\tilde{E})$ -module of vector fields on \tilde{E} . We can write:

$$X = hX + v_1X + v_2X \quad (3.1.)$$

for any $X \in X(\tilde{E})$. It corresponds to the possibility to write any $X \in X(\tilde{E})$ in the following form:

$$X = X^{(h)}(u) \cdot \delta_i + X^{(v_1)}(u) \cdot \delta_i^{(1)} + X^{(v_2)}(u) \cdot \delta_i^{(2)}, \quad u = (x, y^{(1)}, y^{(2)}) \quad (3.2.)$$

We suppose that the manifold M is a paracompact one. In this case \tilde{E} is an orientable Riemannian manifold with respect to Riemannian metric $G^{(\alpha\beta\eta)}$. Thus, general considerations on divergence, gradient and Laplacean done in [5] applies to \tilde{E} . Let us denote by L_X the Lie derivative with respect to $X \in X(\tilde{E})$ and by dV the $3n$ -form on $Osc^2 M$. Then the divergence of $X \in X(\tilde{E})$ is defined by the equation:

$$L_X dV = (div X) dV \quad (3.3.)$$

Let D be a $(\alpha\beta\eta)$ metrical connection. Following the ideas from [1] we obtain:

$$div X = Trace(Y \rightarrow D_Y X + T(X, Y))$$

If we look for the matrix of the linear operator $Y \rightarrow D_Y X + T(X, Y)$ in the adapted basis $\left(\delta_k, \delta_k^{(1)}, \delta_k^{(2)} \right)$ then we obtain:

Theorem 3.1. For any $X = X^{(h)} \cdot \delta_i + X^{(v_1)} \cdot \delta_i^{(1)} + X^{(v_2)} \cdot \delta_i^{(2)} \in X(\tilde{E})$ we have the following formula:

$$\operatorname{div} X = \operatorname{div}_h X + \operatorname{div}_{v_1} X + \operatorname{div}_{v_2} X \quad (3.4)$$

where:

$$\begin{aligned} \operatorname{div}_h X &= X^i_{|i} + X^{(h)k} \left(T^i_{ik} - P^i_{ki} - Q^i_{ki} \right) \\ \operatorname{div}_{v_1} X &= X^i_{||i} + X^{(v_1)k} \left(P^i_{ik} + S^i_{ik} - V^i_{ki} \right) \\ \operatorname{div}_{v_2} X &= X^i_{||i} + X^{(v_2)k} \left(Q^i_{ik} + V^i_{ik} + U^i_{ik} \right) \end{aligned} \quad (3.5)$$

The natural splitting $TOsc^2 M = NOsc^2 M \oplus V_1Osc^2 M \oplus V_2Osc^2 M$ produced by N determines the three partial gradient operators:

$$\begin{aligned} \operatorname{grad}_h f &= g^{ij}(\delta_i f) \cdot \delta_j \\ \operatorname{grad}_{v_1} f &= h^{ij}(\delta_i^{(1)} f) \cdot \delta_j^{(1)} \\ \operatorname{grad}_{v_2} f &= l^{ij}(\delta_i^{(2)} f) \cdot \delta_j^{(2)} \end{aligned} \quad (3.6)$$

and hence we state:

Definition 3.1. Let $f \in F(\tilde{E})$. The vector field

$$\operatorname{grad} f = \operatorname{grad}_h f + \operatorname{grad}_{v_1} f + \operatorname{grad}_{v_2} f \quad (3.7)$$

is called the global gradient with respect to $G^{(\alpha\beta\eta)}$.

Definition 3.2. The scalar field $\Delta f = \operatorname{div}(\operatorname{grad} f)$ is called the Laplacean of the function f with respect to the metrical structure $G^{(\alpha\beta\eta)}$ and the $(\alpha\beta\eta)$ metrical connection D .

For a metrical connection D we can write:

$$\Delta f = \Delta_h f + \Delta_{V_1} f + \Delta_{V_2} f \quad (3.8)$$

where:

$$\begin{aligned} \Delta_h f &= (g^{ij}(\delta_j f))_{;i} + (g^{kp}(\delta_p f)) \cdot \left(T_{ik}^{(1)} - P_{ki}^{(2)} - Q_{ki}^{(3)} \right) \\ \Delta_{V_1} f &= \left(h^{ij}(\delta_j f) \right)_{||i} + \left(h^{kp}(\delta_p f) \right) \cdot \left(P_{ik}^{(1)} + S_{ik}^{(2)} - V_{ki}^{(3)} \right) \\ \Delta_{V_2} f &= \left(l^{ij}(\delta_j f) \right)_{||i} + \left(l^{kp}(\delta_p f) \right) \cdot \left(Q_{ik}^{(1)} + V_{ik}^{(2)} + U_{ik}^{(3)} \right) \end{aligned} \quad (3.9.)$$

Particular cases:

A. Let $G^{(\alpha\beta\eta)}$ be the metrical structure given by (2.1.) and D be the canonical $(\alpha\beta\eta)$ metrical connection given by (2.5.) - (2.10.). In this case we have:

Theorem 3.2. *The following relations hold:*

$$\begin{cases} \operatorname{div} \left(\Gamma^{(1)} \right) = 0, & \operatorname{div} \left(\Gamma^{(2)} \right) = n \left[1 + \left(\frac{\alpha'}{\alpha} + \frac{\beta'}{\beta} + \frac{\eta'}{\eta} \right) \cdot F^2 \right] \\ \Delta E^{(1)} = \frac{2n}{\beta} \left[1 + \left(\frac{\alpha'}{\alpha} + \frac{\beta'}{\beta} + \frac{\eta'}{\eta} \right) \cdot F^2 \right] - \frac{4F^2 \beta'}{\beta^2} \end{cases} \quad (3.10.)$$

where $\Gamma^{(1)}$ and $\Gamma^{(2)}$ are Liouville vector fields given by:

$$\Gamma^{(1)} = y^{(1)i} \cdot \frac{\partial}{\partial y^{(2)i}}, \quad \Gamma^{(2)} = y^{(1)i} \cdot \frac{\partial}{\partial y^{(2)i}} + 2y^{(2)i} \cdot \frac{\partial}{\partial y^{(2)i}} \quad (3.11.)$$

and:

$$E^{(\Phi)} = \Phi \cdot F^2, \quad \Phi = \Phi(F^2) \quad (3.12.)$$

is so called the Φ -energy of first order.

B. If we consider $G^{(0)}$ the homogeneous lift (2.3.) and D the corresponding canonical metrical connection, then we obtain:

Theorem 3.3.

a) *The following relations hold:*

$$\begin{cases} \operatorname{div} \left(\Gamma^{(1)} \right) = 0, & \operatorname{div} \left(\Gamma^{(2)} \right) = -2n \\ \Delta E^{(\Phi)} = 4F^2 \cdot [\Phi'' \cdot F^4 + (3-n) \cdot F^2 \cdot \Phi' + (1-n) \cdot \Phi] \end{cases} \quad (3.13.)$$

b) $E^{(\ast)}$ is harmonic if and only if:

$$\Phi(t) = c_1 \cdot \frac{1}{t} + c_2 \cdot t^{n-1}, \quad c_1, c_2 \in R \quad (3.14.)$$

Aknowledgements. A version of this paper was presented at the Centennial Gh. Vrănceanu, 1st - 4th July, Romanian Academy and University of Bucharest, Romania.

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Received: 01.08.2001

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