

$$(8) \quad \frac{m_1 V^2}{r_1} = \dots$$

**SCHWARZSCHILD'S METRIC GENERATED BY A BODY OF MASS m_1
 PERTURBED BY A MOVING BODY OF MASS m_2**

István Huba Atilla SASS

$$(9) \quad \dots \left(\frac{2m_1}{r_1} - 1 \right) \dots \left(\frac{2m_2}{r_2} - 1 \right) = \dots$$

Abstract. This paper approaches the translation of the coordinate axes of coordinates having the origin O' in the centre of a spherical body of mass m_2 in the reference point with the origin O of a spherical body of mass m_1 , the straight line OO' being the support of the axes $O'x'$ and Ox . This translation being necessary for studying the influence of the body with mass m_2 on that with mass m_1 . We assume that the body of mass m_2 moves compared to the body of mass m_1 at a radial speed V .

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1. Notations.

Let us consider two spherical bodies of masses m_1 and m_2 , respectively, relatively for-off so that their shapes could be considered spherical.

In this case a references point can be chosen for both bodies, so that mass m_1 should be in the origin O of the system of axes $Oxyzt$ and mass m_2 should be in the origin O' of the system of axes $O'x'y'z'$.

We choose the axes Ox and $O'x'$ on the straight line OO' .

Schwarzschild's metrics, associated with the body of mass m_1 situated in the origin of the references point $Oxyzt$.

$$ds_1^2 = \left(1 - \frac{2m_1}{r}\right) dt^2 - \left(1 - \frac{2m_1}{r}\right)^{-1} dr^2 - r^2 d\theta^2 - r^2 \sin^2 \theta d\phi^2 = g_{ij} dx^i dx^j \quad (1)$$

where

$$(2) \quad \left. \begin{aligned} x^0 = t = ct_{ph}; \quad x^1 = r = r_{ph}; \quad x^2 = \theta = \theta_{ph}; \quad x^3 = \phi = \phi_{ph} \end{aligned} \right\}$$

are the generalized and geometrized spherical coordinates. The geometrized time and mass have length dimensions (in meters).

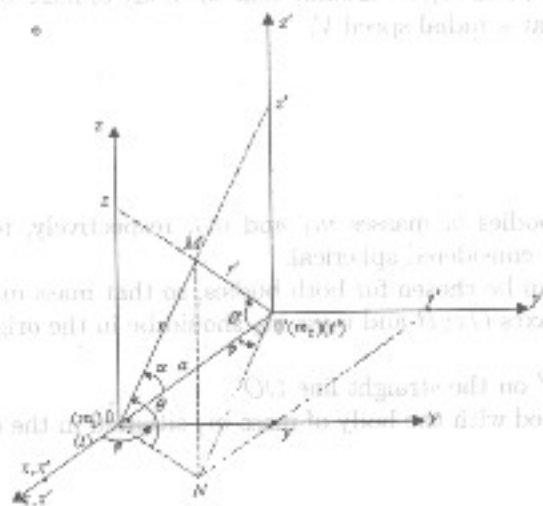
$$m = \frac{GM_{ph}}{c^2} \quad (3)$$

where index *ph* showing that the respective magnitude is measured in physical units and *c* and *G* represent the speed of the light and the universal constant of gravity, respectively.

To mass m_2 situated in the origin O' of the axes of coordinates $O'x'y'z't'$ corresponds the metrics

$$ds_2^2 = \left(1 - \frac{2m_2}{r^1}\right) dt'^2 - \left(1 - \frac{2m_2}{r^1}\right)^{-1} dr'^2 - r'^2 d\theta'^2 - r'^2 \sin^2 \theta' d\phi'^2 \quad (4)$$

In both cases, we will consider the meter as a light unit. A point M from the spacetime Universe has the Cartesian coordinates (x, y, z, t) measured from O and (x', y', z', t') measured from O' . We have noted the radial distances by $OM \equiv r$; $O'M \equiv r'$; $OO' \equiv a$ (5)



(1) (see figure 1) The relation between the Cartesian coordinates and the spherical ones is given by the formulas:

$$\begin{cases} x = r \cos \theta \cos \phi \\ y = r \cos \theta \sin \phi \\ z = r \sin \theta \end{cases} \quad \begin{cases} x' = r' \cos \theta' \cos \phi' \\ y' = r' \cos \theta' \sin \phi' \\ z' = r' \sin \theta' \end{cases} \quad (6)$$

On the other hand

$$(7) \quad \begin{cases} x' = a + x \\ y' = y \\ z' = z \end{cases} \quad (7)$$

In the triangle MMO' we have

$$(8) \quad r'^2 = a^2 + r^2 - 2ar \cos \alpha$$

But

$$(9) \quad r \cos(180^\circ - \alpha) = -r \cos \alpha = x,$$

from (8) and (9) it results that $r'^2 = a^2 + r^2 + 2ax$, and replacing x from first formula (6) we get

$$(10) \quad r' = E = \sqrt{a^2 + r^2 + 2ar \cos \theta \cos \phi}$$

An event occurring in the point M at the moment $t_{M,a}$ measured after the clock in M' will be seen in O after the arrival of the luminous signal at the speed of the light c at the moment $t_{pa} + \frac{r}{c}$, and in O' at the moment $t'_{pa} + \frac{r'}{c}$. Passing to generalized units:

$$(11) \quad t_M = t + r = t' + r'$$

and consequently

$$(12) \quad t' = t + r - E$$

As $z = z'$ from (6) and (10), it results

$$(13) \quad \sin \theta' = \frac{r \sin \theta}{E}$$

and

$$(14) \quad \cos \theta' = \frac{F}{E}$$

where

$$(15) \quad F = \sqrt{a^2 + r^2 \cos^2 \theta + 2ar \cos \theta \cos \phi}$$

As $y' = y$, from (6) and (15) also results that

$$(16) \quad \sin \phi' = \frac{r \cos \theta \sin \phi}{F}$$

and

$$(7) \quad \cos \phi' = \frac{a + r \cos \theta \cos \phi}{F} \quad (17)$$

2. The transcription of Schwarzschild's metric associated with mass m_2 in the

coordinates from m_1

(8) Supposing that mass m_2 moves as referred to m_1 , at a radial speed V_{ph} that is

$$(9) \quad da = V dt \quad (18)$$

where

$$(10) \quad V = \frac{V_{ph}}{c} \quad (19)$$

From the formula

$$(11) \quad r' \sin \theta = r \sin \phi \quad (20)$$

and differentiating the formulae (10), (12), (16), (13) and (15), we get

$$(12) \quad dr' = dE = \frac{1}{E} [(a + r \cos \theta \cos \phi) V dt + (r + a \cos \theta \cos \phi) dr - ar \sin \theta \cos \phi d\theta - ar \cos \theta \sin \phi d\phi] \quad (21)$$

$$(13) \quad dt' = dt + dr - dE \quad (22)$$

$$(14) \quad d\theta' = \frac{1}{F} \left(\sin \theta dr + r \cos \theta d\theta - r \sin \theta \frac{dE}{E} \right) \quad (23)$$

$$(15) \quad d\phi' = \frac{1}{a + r \cos \theta \cos \phi} \left(\cos \theta \sin \phi dr - r \sin \theta \sin \phi d\theta + r \cos \theta \cos \phi d\phi - r \cos \theta \sin \phi \frac{dF}{F} \right) \quad (24)$$

$$(16) \quad dF = \frac{1}{F} [(a + r \cos \theta \cos \phi) V dt + (r \cos \theta + a \cos \phi) \cos \theta dr - (r + a) r \sin \theta \cos \phi d\theta - ar \cos \theta \sin \phi d\phi] \quad (25)$$

After the corresponding replacements in metrics (4), it is demonstrated the following theorem

Theorem. *The Schwarzschild's metric generated by a body with mass m_1 is perturbed by moving body with mass m_2 with radial speed V situated at a distance a , with the supplementary metric*

$$ds_2^2 = \frac{1}{E^2} \left(1 - \frac{2m_2}{E} \right) A^2 - \frac{1}{E^2} \left(1 - \frac{2m_2}{E} \right)^{-1} B^2 - \frac{1}{E^2 F^2} C^2 - \frac{r^2 \sin^2 \theta}{F^2 (a - r \cos \theta \cos \phi)^2} D^2 \quad (26)$$

where

$$A = [E - (a + r \cos \theta \cos \phi)V dt + [E - (r + a \cos \theta \cos \phi)] dr + ar \sin \theta \cos \phi d\theta + ar \cos \theta \sin \phi d\phi] \quad (27)$$

$$B = (a + r \cos \theta \cos \phi)V dt + (r + a \cos \theta \cos \phi) dr - ar \sin \theta \cos \phi d\theta - ar \cos \theta \sin \phi d\phi \quad (28)$$

$$C = (a + r \cos \theta \cos \phi)V r \sin \theta dt - a \sin \theta (a + r \cos \theta \cos \phi) dr - r [(a^2 + r^2) \cos \theta + ar \cos \phi (1 + \cos^2 \theta)] d\theta - ar^2 \sin \theta \cos \theta \sin \phi d\phi \quad (29)$$

$$D = (a + r \cos \theta \cos \phi)V r \cos \theta \sin \phi dt - a \cos \theta \sin \phi (a + r \cos \theta \cos \phi) dr + r \sin \theta \sin \phi [a^2 + r^2 \cos^2 \theta + r(a - r) \cos \theta + \cos \theta \cos \phi] d\theta - r \cos \theta [ar \cos \theta (1 + \cos^2 \phi) + (a^2 + r^2 \cos^2 \theta) \cos \phi] d\phi \quad (30)$$

□

Metrics (26) shows us the modifications that have been brought about by the body of mass m_2 of the space-time seen from the body of mass m_1 .

The body of mass m_2 moves of a radial speed v , compared to the body of mass m_1 . In our case, metrics (26) differs from the corresponding metrics (16) from the previous paper for the static case, only through the temporary components h_{ij} which also include O indices.

REFERENCES

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Department of Mathematics and Computer Science
 North University of Baia Mare, Str. Victoriei nr. 76
 4800 Baia Mare, ROMANIA

$$\begin{aligned}
 (1) \quad & -\partial_t(\partial_x \rho + \partial_x p) + \partial_x(\rho v) + \partial_x(p v) = 0 \\
 (2) \quad & -\partial_t(\partial_x \rho v) + \partial_x(\rho v^2) + \partial_x(\rho v p) = 0 \\
 (3) \quad & -\partial_t(\partial_x \rho v^2) + \partial_x(\rho v^3) + \partial_x(\rho v^2 p) = 0 \\
 (4) \quad & -\partial_t(\partial_x \rho v^3) + \partial_x(\rho v^4) + \partial_x(\rho v^3 p) = 0 \\
 (5) \quad & -\partial_t(\partial_x \rho v^4) + \partial_x(\rho v^5) + \partial_x(\rho v^4 p) = 0 \\
 (6) \quad & -\partial_t(\partial_x \rho v^5) + \partial_x(\rho v^6) + \partial_x(\rho v^5 p) = 0
 \end{aligned}$$

The above system of equations is solved numerically using the method of characteristics. The initial conditions are given by the Rankine-Hugoniot conditions across the shock front. The numerical results show that the shock front propagates with a constant velocity and the pressure and density profiles behind the shock front are self-similar. The numerical results are compared with the analytical solutions obtained in [1].