

# Stability results of a suspension-bridge with nonlinear damping modulated by a time dependent coefficient

MOHAMMAD M. AL-GHARABLI<sup>1</sup> and SALIM A. MESSAOUDI<sup>2</sup>

**ABSTRACT.** The main goal of this work is to investigate the following weakly damped nonlinear suspension-bridge equation

$$u_{tt}(x, y, t) + \Delta^2 u(x, y, t) + \alpha(t)g(u_t) = 0,$$

and establish explicit and general decay results for the energy of solutions of the problem. Our decay results depend on the functions  $\alpha$  and  $g$  and obtained without any restriction growth assumption on  $g$  at the origin. The multiplier method, the properties of the convex and the dual of the convex functions, Jensen's inequality and the generalized Young inequality are used to establish the stability results.

## 1. INTRODUCTION

The importance of bridges is undeniable and their presence in the human daily life goes back in history for a long time. Though their importance, bridges have brought some challenges, such as collapse and instability due to nature hazards such as winds, earthquakes, ... etc. To overcome these difficulties, engineers and scientists have made efforts to find the best designs and models possible. Many mathematical models have appeared since the collapse of Tacoma Narrows Bridge. Motivated by the wonderful book of Rocard [17], where it was pointed out that the correct way to model a suspension bridge is through a thin plate, Ferrero-Gazzola [8] introduced the following hyperbolic problem:

$$(1.1) \quad \begin{cases} u_{tt}(x, y, t) + \eta u_t + \Delta^2 u(x, y, t) + h(x, y, u) = f, & \text{in } \Omega \times \mathbb{R}^+, \\ u(0, y, t) = u_{xx}(0, y, t) = u(\pi, y, t) = u_{xx}(\pi, y, t) = 0, & (y, t) \in (-l, l) \times \mathbb{R}^+, \\ u_{yy}(x, \pm l, t) + \sigma u_{xx}(x, \pm l, t) = 0, & (x, t) \in (0, \pi) \times \mathbb{R}^+, \\ u_{yyy}(x, \pm l, t) + (2 - \sigma)u_{xyy}(x, \pm l, t) = 0, & (x, t) \in (0, \pi) \times \mathbb{R}^+, \\ u(x, y, 0) = u_0(x, y), \quad u_t(x, y, 0) = u_1(x, y), & \text{in } \Omega \times \mathbb{R}^+, \end{cases}$$

where  $\Omega = (0, \pi) \times (-l, l)$  is a planar rectangular plate,  $\sigma$  is the well-known Poisson ratio,  $\eta$  is the damping coefficient,  $h$  is the nonlinear restoring force of the hangers and  $f$  is an external force. After the appearance of the above model, many mathematician showed interest in investigating variants of it, using different kinds of damping in the aim to obtain stability of the bridge modelled though the above problem. Wang [18] considered the equation

$$u_{tt} + \Delta^2 u + \nu u_t = |u|^{p-2}u,$$

together with the above initial and boundary conditions. After showing the uniqueness and existence of local solutions, he gave sufficient conditions for global existence and finite-time blow-up of solutions. Messaoudi and Mukaiawa [16] studied the above problem, where the linear frictional damping was replaced by a nonlinear frictional damping and established the existence of a global weak solution and proved exponential and polynomial stability results. Recently, Audu et al. [4] considered a plate equation as

Received: 24.06.2022. In revised form: 10.01.2023. Accepted: 17.01.2023

2020 Mathematics Subject Classification. 35B35, 35L55, 75D05, 74D10, 93D20.

Key words and phrases. Suspension-bridge, Plate equation, General decay, Nonlinear frictional damping.

Corresponding author: Mohammad M. Al-Gharabli; mahfouz@kfupm.edu.sa

a model for a suspension bridge with a general nonlinear internal feedback and time-varying weight. Under some conditions on the feedback and the coefficient functions, they established a general decay estimate. More results in this direction can be found in [1, 15, 6, 5, 9, 10, 14, 7].

Our aim in this work is to investigate the following plate problem

$$(1.2) \quad \begin{cases} u_{tt}(x, y, t) + \Delta^2 u(x, y, t) + \alpha(t)g(u_t) = 0, & \text{in } \Omega \times (0, T), \\ u(0, y, t) = u_{xx}(0, y, t) = u(\pi, y, t) = u_{xx}(\pi, y, t) = 0, & (y, t) \in (-d, d) \times (0, T), \\ u_{yy}(x, \pm d, t) + \sigma u_{xx}(x, \pm d, t) = 0, & (x, t) \in (0, \pi) \times (0, T), \\ u_{yyy}(x, \pm d, t) + (2 - \sigma)u_{xxy}(x, \pm d, t) = 0, & (x, t) \in (0, \pi) \times (0, T), \\ u(x, y, 0) = u_0(x, y), \quad u_t(x, y, 0) = u_1(x, y), & \text{in } \Omega \times (0, T), \end{cases}$$

where  $\Omega = (0, \pi) \times (-d, d)$ ,  $d \ll \pi$ ,  $g : \mathbb{R} \rightarrow \mathbb{R}$  and  $\alpha : [0, +\infty) \rightarrow (0, +\infty)$  is a nonincreasing differentiable function,  $u$  is the vertical displacement of the bridge and  $\sigma$  is the Poisson ratio. For metals its value lies around 0.3, while for concrete it is between 0.1 and 0.2. For this reason, we shall assume that  $0 < \sigma < \frac{1}{2}$ .

This is a weakly damped nonlinear suspension-bridge problem, in which the damping is modulated by a time dependent coefficient  $\alpha(t)$ . We establish an explicit and general decay result, depending on  $g$  and  $\alpha$ , for which the optimal exponential and polynomial decay rate estimates are only special cases. The proof is based on the multiplier method and makes use of some properties of convex functions including the use of the general Young inequality and Jensen’s inequality. The proof of the current result is easier than the one in [4]. Moreover, this result gives a better rate of decay (see Remark 3.3 below).

The paper is organized as follows. In section 2, we present some notations and material needed for our work. The statement and the proof of our main result will be given in section 3.

## 2. PRELIMINARIES

In this section, we present some material needed in the proof of our results. First, we introduce the space

$$(2.3) \quad H_*^2(\Omega) = \{w \in H^2(\Omega) : w = 0 \text{ on } \{0, \pi\} \times (-d, d)\},$$

together with the inner product

$$(2.4) \quad (u, v)_{H_*^2} = \int_{\Omega} (\Delta u \Delta v + (1 - \sigma)(2u_{xy}v_{xy} - u_{xx}v_{yy} - u_{yy}v_{xx})) \, dx \, dy.$$

It is well known that  $(H_*^2(\Omega), (\cdot, \cdot)_{H_*^2})$  is a Hilbert space, and the norm  $\|\cdot\|_{H_*^2}^2$  is equivalent to the usual  $H^2$ , see [8]. We also let

$$(2.5) \quad \mathcal{H}(\Omega) := \text{The dual space of } H_*^2(\Omega).$$

Throughout this paper,  $c$  is used to denote a generic positive constant.

**Lemma 2.1.** [18] *Let  $u \in H_*^2(\Omega)$  and assume that  $1 \leq p < \infty$ , then, there exists a positive constant  $C_e = C_e(\Omega, p) > 0$  such that*

$$\|u\|_p^p \leq C_e \|u\|_{H_*^2(\Omega)}^p.$$

**Lemma 2.2.** (Jensen’s inequality) *Let  $G : [a, b] \rightarrow \mathbb{R}$  be a convex function. Assume that the functions  $f : (0, L) \rightarrow [a, b]$  and  $r : (0, L) \rightarrow \mathbb{R}$  are integrable such that  $r(x) \geq 0$ , for any  $x \in (0, L)$  and  $\int_0^L r(x)dx = k > 0$ . Then,*

$$(2.6) \quad G \left( \frac{1}{k} \int_0^L f(x)r(x)dx \right) \leq \frac{1}{k} \int_0^L G(f(x))r(x)dx.$$

We consider the following hypotheses

(H1)  $g : \mathbb{R} \rightarrow \mathbb{R}$  is nondecreasing  $C^1$  function such that there exists a  $C^2$  convex and increasing function  $G : \mathbb{R}^+ \rightarrow \mathbb{R}^+$  satisfying  $G(0) = 0$  and  $G'(0) = 0$  or  $G$  is linear on  $[0, \varepsilon]$  such that

$$(2.7) \quad \begin{aligned} c_1|s| \leq |g(s)| \leq c_2|s|, & \text{ if } |s| \geq \varepsilon, \\ |s|^2 + g^2(s) \leq G^{-1}(sg(s)), & \text{ if } |s| \leq \varepsilon, \end{aligned}$$

where  $\varepsilon, c_1, c_2$  are positive constant.

(H2)  $\alpha : [0, +\infty) \rightarrow (0, +\infty)$  is a nonincreasing differentiable function.

**Remark 2.1.** Hypothesis (H1) implies that  $sg(s) > 0$ , for all  $s \neq 0$ .

**Remark 2.2.** If  $g$  satisfies

$$(2.8) \quad \begin{aligned} g_0(|s|) \leq |g(s)| \leq g_0^{-1}(|s|), & \quad |s| \leq \varepsilon \\ c_1|s| \leq |g(s)| \leq c_2|s|, & \quad |s| \geq \varepsilon \end{aligned}$$

for some strictly increasing function  $g_0 \in C^1([0, +\infty))$ , with  $g_0(0) = 0$ , and positive constants  $c_1, c_2, \varepsilon$  and the function  $G(s) = \sqrt{\frac{s}{2}}g_0\left(\frac{s}{2}\right)$ , is strictly convex  $C^2$  function on  $(0, \varepsilon]$  when  $g_0$  is nonlinear, then (H1) is satisfied. This kind of hypotheses, where (H1) is weaker, was considered by Liu and Zuazua [13] and Alabau-Boussouira [2].

For completeness we state the following existence and regularity result whose proof can be established similarly to that in [12] and [16].

**Theorem 2.1.** Let  $(u_0, u_1) \in H_*^2(\Omega) \times L^2(\Omega)$ . Assume that (H1) and (H2) hold. Then problem (1.2) has a unique weak global solution

$$u \in C([0, T], H_*^2(\Omega)), \quad u_t \in C([0, T], L^2(\Omega)).$$

The energy functional associated with problem (1.2) is

$$(2.9) \quad E(t) = \frac{1}{2} \left( \|u_t\|_2^2 + \|u\|_{H_*^2(\Omega)}^2 \right).$$

Direct differentiation of (2.9), using (1.2), leads to

$$(2.10) \quad E'(t) = -\alpha(t) \int_{\Omega} u_t g(u_t) dx \leq 0.$$

### 3. MAIN RESULT

In this section, we state and prove our main result. We first start with the following lemma.

**Lemma 3.3.** The functional

$$L(t) = NE(t) + \int_{\Omega} uu_t dx$$

satisfies, along the solution of (1.2),

$$(3.11) \quad L'(t) \leq -c_1 E(t) + c_2 \int_{\Omega} (u_t^2 + g^2(u_t)) dx$$

and

$$(3.12) \quad L \sim E,$$

where  $c_1$  and  $c_2$  are positive constants.

*Proof.* Differentiate  $L(t)$ , using (1.2), (2.9) and (2.10), to get

$$(3.13) \quad L'(t) = NE'(t) + \|u_t\|_2^2 - \|u\|_{H_*^2(\Omega)}^2 - \alpha(t) \int_{\Omega} ug(u_t)dx.$$

The use of Young's inequality and the embedding lemma 2.1 gives

$$L'(t) \leq NE'(t) + \|u_t\|_2^2 - (1 - c\delta)\|u\|_{H_*^2(\Omega)}^2 + c(\delta) \int_{\Omega} g^2(u_t)dx.$$

Selecting  $\delta$  small enough and  $N$  large enough, we easily obtain (3.11) and (3.12). □

We are now ready to state and prove our main result.

**Theorem 3.2.** *Let  $(u_0, u_1) \in H_*^2(\Omega) \times L^2(\Omega)$ . Assume that (A1)- (A2) hold. Then there exist positive constants  $\varepsilon_0, d_1, d_2, k_1, k_2$  and  $k_3$  such that*

$$(3.14) \quad E(t) \leq d_1 e^{-d_2 \int_0^t \alpha(s)ds}, \text{ if } G \text{ is linear}$$

and

$$(3.15) \quad E(t) \leq k_1 W_1^{-1}(k_2 \int_0^t \alpha(s)ds + k_3), \text{ otherwise,}$$

where

$$W_1(\tau) = \int_{\tau}^1 \frac{1}{W_2(s)} \text{ and } W_2(t) = tG'(\varepsilon_0 t).$$

*Proof.* We multiply (3.11) by  $\alpha(t)$  to get

$$(3.16) \quad \alpha(t)L'(t) \leq -c_1\alpha(t)E(t) + c_2\alpha(t) \int_{\Omega} (u_t^2 + g^2(u_t)) dx.$$

**Case 1.**  $G$  is linear on  $[0, \varepsilon]$ . Then, using (A1) and (2.10), estimate (3.16) becomes

$$\alpha(t)L'(t) \leq -c_1\alpha(t)E(t) + c\alpha(t) \int_{\Omega} u_t g(u_t)dx = -c_1\alpha(t)E(t) - cE'(t),$$

which gives,

$$(\alpha L + cE)' \leq -c_1\alpha(t)E(t).$$

Hence, using the fact that  $\alpha L + cE \sim E$ , we get (3.14).

**Case 2.**  $G$  is nonlinear on  $[0, \varepsilon]$ .

Let  $0 < \varepsilon_1 \leq \varepsilon$  such that

$$(3.17) \quad sg(s) \leq \min\{\varepsilon, G(\varepsilon)\} \text{ for all } |s| \leq \varepsilon_1.$$

Recalling (H1) and Remark 2.2, we have, for  $\varepsilon_1 \leq |s| \leq \varepsilon$ ,

$$|g(s)| \leq \frac{g_0^{-1}(|s|)}{|s|} |s| \leq \frac{g_0^{-1}(|\varepsilon|)}{|\varepsilon_1|} |s| = c'_2 |s|$$

and

$$|g(s)| \geq \frac{g_0(|s|)}{|s|} |s| \geq \frac{g_0(|\varepsilon_1|)}{|\varepsilon|} |s| = c'_1 |s|.$$

Therefore, we deduce that

$$(3.18) \quad \begin{cases} s^2 + g^2(s) \leq G^{-1}(sg(s)) & \text{for all } |s| \leq \varepsilon_1 \\ c'_1|s| \leq |g(s)| \leq c'_2|s| & \text{for all } |s| \geq \varepsilon_1. \end{cases}$$

To estimate the last integral in (3.16), we use the following partition which was first introduced by Komornik [11]:

$$\Omega_1 = \{x \in \Omega : |u_t| \leq \varepsilon_1\}, \quad \Omega_2 = \{x \in \Omega : |u_t| > \varepsilon_1\}.$$

With

$$J(t) := \frac{1}{|\Omega_1|} \int_{\Omega_1} u_t g(u_t) dx,$$

Jensen's inequality gives (note that  $G^{-1}$  is concave)

$$(3.19) \quad G^{-1}(J(t)) \geq c \int_{\Omega_1} G^{-1}(u_t g(u_t)) dx.$$

Thus, combining (2.10), (3.18) and (3.19), we arrive at

$$(3.20) \quad \begin{aligned} \alpha(t) \int_{\Omega} (u_t^2 + g^2(u_t)) dx &\leq \alpha(t) \int_{\Omega_1} G^{-1}(u_t g(u_t)) dx + \alpha(t) \int_{\Omega_2} (u_t^2 + g^2(u_t)) dx \\ &\leq c\alpha(t)G^{-1}(J(t)) + c\alpha(t) \int_{\Omega_2} u_t g(u_t) dx \\ &\leq c\alpha(t)G^{-1}(J(t)) - cE'(t). \end{aligned}$$

A combination of (3.16) and (3.20) yields

$$(3.21) \quad L'_1(t) \leq -c_1\alpha(t)E(t) + c\alpha(t)G^{-1}(J(t)), \quad \forall t \in \mathbb{R}^+,$$

where  $L_1 = \alpha L + cE$ , which is clearly equivalent to  $E$ . Now, for  $\varepsilon_0 < \varepsilon$  and  $c_0 > 0$ , let

$$L_2(t) := G' \left( \varepsilon_0 \frac{E(t)}{E(0)} \right) L_1(t) + c_0 E(t),$$

provided that  $E(0) > 0$ ; otherwise  $E(t) = 0, \forall t \in \mathbb{R}^+$ , hence the theorem is verified. By using (3.21) and exploiting the properties of  $E$  and  $G$ , we conclude that  $L_2$  satisfies

$$(3.22) \quad \begin{aligned} L'_2(t) &= \varepsilon_0 \frac{E'(t)}{E(0)} G'' \left( \varepsilon_0 \frac{E(t)}{E(0)} \right) L_1(t) + G' \left( \varepsilon_0 \frac{E(t)}{E(0)} \right) L'_1(t) + c_0 E'(t) \\ &\leq -c_1\alpha(t)E(t)G' \left( \varepsilon_0 \frac{E(t)}{E(0)} \right) + c\alpha(t)G^{-1}(J(t))G' \left( \varepsilon_0 \frac{E(t)}{E(0)} \right) \\ &\quad + c_0 E'(t). \end{aligned}$$

Let  $G^*$  be the convex conjugate of  $G$  in the sense of Young (see [3], p.61-64), then, for  $s \in (0, G'(\varepsilon)]$ ,

$$(3.23) \quad G^*(s) = s(G')^{-1}(s) - G[(G')^{-1}(s)] \leq s(G')^{-1}(s).$$

Using the general Young inequality

$$AB \leq G^*(A) + G(B), \quad \text{if } A \in (0, G'(\varepsilon)], \quad B \in (0, \varepsilon],$$

for

$$A = G' \left( \varepsilon_0 \frac{E(t)}{E(0)} \right) \quad \text{and} \quad B = G^{-1}(J(t)),$$

we get

$$\begin{aligned} L'_3(t) &\leq -c_1\alpha(t)E(t)G' \left( \varepsilon_0 \frac{E(t)}{E(0)} \right) + c\alpha(t)\varepsilon_0 \frac{E(t)}{E(0)}G' \left( \varepsilon_0 \frac{E(t)}{E(0)} \right) - cE'(t) \\ &\quad + c_0E'(t) \\ &= -(c_1E(0) - C\varepsilon_0)\alpha(t) \frac{E(t)}{E(0)}G' \left( \varepsilon_0 \frac{E(t)}{E(0)} \right) - (c - c_0)E'(t). \end{aligned}$$

Consequently, by picking  $\varepsilon_0$  small enough so that  $c_1E(0) - C\varepsilon_0 > 0$  and  $c_0$  large such that  $c - c_0 < 0$ , we obtain, for all  $t \in \mathbb{R}^+$ ,

$$(3.24) \quad \alpha_1L_2(t) \leq E(t) \leq \alpha_2L_2(t),$$

for some  $\alpha_1, \alpha_2 > 0$  and for some  $k > 0$ ,

$$(3.25) \quad L'_2(t) \leq -k\alpha(t) \frac{E(t)}{E(0)}G' \left( \varepsilon_0 \frac{E(t)}{E(0)} \right) = -k\alpha(t)W_2 \left( \frac{E(t)}{E(0)} \right),$$

where

$$W_2(s) = sG'(\varepsilon_0s).$$

Thus, with  $R(t) = \frac{\alpha_1L_2(t)}{E(0)}$  and using (3.24) and (3.25), we have

$$(3.26) \quad R(t) \sim E(t)$$

and, for some  $k_2 > 0$ ,

$$(3.27) \quad R'(t) \leq -k_2\alpha(t)W_2(R(t)).$$

Inequality (3.27) implies that  $(W_1(R))' \geq k_2\alpha(t)$ , where

$$W_1(\tau) = \int_{\tau}^1 \frac{1}{W_2(s)} ds \quad \text{for } \tau \in (0, 1].$$

So, by integrating over  $[0, t]$ , we get

$$(3.28) \quad R(t) \leq W_1^{-1} \left( k_2 \int_0^t \alpha(s) ds + k_3 \right), \quad \forall t \in \mathbb{R}^+.$$

Finally, we obtain (3.15) by virtue of (3.26) and (3.28). □

**Example 3.1.** We give some examples to illustrate the energy decay rates given by Theorem 3.2. Here, we assume that  $g$  satisfies (2.8) near the origin with the following various examples for  $g_0$ :

1. If  $g_0(s) = cs^q$  and  $q \geq 1$ , then  $G(s) = cs^{\frac{q+1}{2}}$  satisfies (A1) and, consequently, Theorem (3.2) yields

$$E(t) \leq d_1e^{-d_2 \int_0^t \alpha(s) ds}, \quad \text{if } q = 1,$$

$$E(t) \leq k_1 \left( k_2 \int_0^t \alpha(s) ds + k_3 \right), \quad \text{if } q > 1.$$

2. If  $g_0(s) = e^{-\frac{1}{s}}$ , then (A1) is satisfied for  $G(s) = \sqrt{\frac{s}{2}}e^{-\sqrt{\frac{2}{s}}}$  near zero. Therefore, we get

$$E(t) \leq k_1 \left( \ln \left( k_2 \int_0^t \alpha(s) ds + k_3 \right) \right)^{-2}.$$

**Remark 3.3.** This work improves the work of Audu et al. [4] in several aspects: first, we assume less conditions on the nonlinear feedback, this allows a wider class of functions. Second, our decay result is explicit and clearer than the one obtained in [4]. Third, it is easy to compute the decay rate of the well-known nonlinear feedback.

**Remark 3.4.** Note that the exponential and the polynomial decay results are only special cases.

**Acknowledgment.** The authors would like to express their profound gratitude to King Fahd University of Petroleum and Minerals (KFUPM) for its continuous support. This work is funded by KFUPM under Project #SB201003.

## REFERENCES

- [1] Al-Gwaiz, M.; Benci, V.; Gazzola, F. Bending and stretching energies in a rectangular plate modeling suspension bridges. *Nonlinear Anal.* **106** (2014), 18–34.
- [2] Alabau-Boussouira, F. Convexity and weighted integral inequalities for energy decay rates of nonlinear dissipative hyperbolic systems. *Appl. Math. Comput.* **51** (2005), no. 1, 61–105.
- [3] Arnold, V. *Mathematical methods of classical mechanics*. (Vol. 60) Springer Science & Business Media, (2013).
- [4] Audu, J.; Mukaiawa, S.; Júnior, D. General decay estimate for a two-dimensional plate equation with time-varying feedback and time-varying coefficient. *Results Appl Math.* **12** (2021), Article 100219.
- [5] Cavalcanti, A.; Cavalcanti, M.; Corrêa, W.; Hajje, Z.; Cortés, M.; Asem, R. Uniform decay rates for a suspension bridge with locally distributed nonlinear damping. *J. Franklin Inst.* **357** (2020), no. 4, 2388–2419.
- [6] Cavalcanti, M.; Corrêa, W.; Fukuoka, R.; Hajje, Z. Stabilization of a suspension bridge with locally distributed damping. *Math. Control Signals Syst.* **30** (2018), no. 4, 1–39.
- [7] Ferreira, Jr.; Gazzola, F.; dos Santos, E. Instability of modes in a partially hinged rectangular plate. *J. Differ. Equ.* **261** (2016), no. 11, 6302–6340.
- [8] Ferrero, A.; Gazzola, F. A partially hinged rectangular plate as a model for suspension bridges. *Discrete and continuous dynamical systems* **35** (2015), no. 12, 5879–5908.
- [9] Gazzola, F. *Mathematical models for suspension bridges*. Cham: Springer, (2015).
- [10] Hajje, Z.; Messaoudi, S. Stability of a suspension bridge with structural damping. In *Annales Polonici Mathematici* (Vol. 125, pp. 59–70). Instytut Matematyczny Polskiej Akademii Nauk, (2020).
- [11] Komornik, V. *Exact controllability and stabilization: the multiplier method* (Vol. 36). Elsevier Masson, (1994).
- [12] Lasiecka, I.; Tataru, D. Uniform boundary stabilization of semilinear wave equations with nonlinear boundary damping. *Differential and Integral Equations* **6** (1993), no. 3, 507–533.
- [13] Liu, W.; Zuazua, E. Decay rates for dissipative wave equations. *Ricerche di Matematica* **48** (1999), 61–75.
- [14] Liu, W.; Zhuang, H. Global existence, asymptotic behavior and blow-up of solutions for a suspension bridge equation with nonlinear damping and source terms. *Nonlinear Differential Equations and Applications NoDEA* **24** (2017), no. 6, 1–35.
- [15] Messaoudi, S.; Mukaiawa, S. A suspension bridge problem: existence and stability. In *International Conference on Mathematics and Statistics* (pp. 151–165). Springer, Cham, (2017).
- [16] Messaoudi, S.; Mukaiawa, S. Existence and stability of fourth-order nonlinear plate problem. *Nonautonomous Dynamical Systems* **6** (2019), no. 1, 81–98.
- [17] Rocard, Y. *Dynamic instability: automobiles, aircraft, suspension bridges*. C. Lockwood, (1957).
- [18] Wang, Y. Finite time blow-up and global solutions for fourth order damped wave equations. *J. Math. Anal. Appl.* **418** (2014), no. 2, 713–733.

<sup>1</sup>THE INTERDISCIPLINARY RESEARCH CENTER IN CONSTRUCTION AND BUILDING MATERIALS  
 KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS  
 DHAHRAN 31261, SAUDI ARABIA  
 Email address: mahfouz@kfupm.edu.sa

<sup>2</sup>DEPARTMENT OF MATHEMATICS  
 UNIVERSITY OF SHARJAH  
 P. O. BOX, 27272, SHARJAH. UAE  
 Email address: smessaoudi@sharjah.ac.ae